

# Tutorial 6 Solutions

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1. \*Show that

$$\psi(x) = A \sin(kx) + B \cos(kx)$$

and

$$\psi(x) = C e^{ikx} + D e^{-ikx}$$

are equivalent solutions of TISE of a free particle. A, B, C and D can be complex numbers.

Sol<sup>n</sup>

$$\begin{aligned} \psi(x) = A \sin kx + B \cos kx &= A \left( \frac{e^{ikx} - e^{-ikx}}{2i} \right) + B \left( \frac{e^{ikx} + e^{-ikx}}{2} \right) \\ &= e^{ikx} \left( \frac{A}{2i} + \frac{B}{2} \right) + e^{-ikx} \left( \frac{-A}{2i} + \frac{B}{2} \right) \end{aligned}$$

2. Show that

$$\Psi(x, t) = A \sin(kx - \omega t) + B \cos(kx - \omega t)$$

does not obey the time-dependant Schroedinger's equation for a free particle.

Sol<sup>n</sup> TDSE

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \Psi(x, t)}{\partial x^2} = i\hbar \frac{\partial \Psi(x, t)}{\partial t}$$

$$-\frac{\hbar^2 k^2}{2m} \Psi(x, t) = -i\hbar \omega A \cos(kx - \omega t) + i\hbar \omega B \sin(kx - \omega t)$$

Clearly  $\Psi(x, t)$  is not an eigenfunction of  $(-i\hbar \frac{\partial}{\partial t})$ . Hence it isn't a valid sol<sup>n</sup> of TDSE.

4. \* A free proton has a wave function given by

$$\Psi(x, t) = A e^{i(5.02 \cdot 10^{11} x - 8.00 \cdot 10^{15} t)} = A e^{i(kx - \omega t)}$$

The coefficient of  $x$  is inverse meters, and the coefficient of  $t$  is inverse seconds. Find its momentum and energy.

Sol<sup>n</sup>  $\Psi(x, t) = A e^{i\alpha x} \cdot e^{-i\beta t} = \underbrace{\Psi(x)}_{\phi(x)} \cdot \underbrace{T(t)}_{T(t)}$

$\hookrightarrow$  eigen function of both  $\hat{p} \left( -i\hbar \frac{\partial}{\partial x} \right)$  &  $\hat{H} \left( i\hbar \frac{\partial}{\partial t} \right)$

$$\hat{p} \Psi(x, t) = -i\hbar \frac{\partial}{\partial x} (A e^{i\alpha x} \cdot e^{-i\beta t}) = \hbar \alpha (A e^{i\alpha x} \cdot e^{-i\beta t}) = \hbar \alpha \Psi(x, t)$$

$$\boxed{p = \hbar \alpha}$$

$$\boxed{E = \hbar \beta}$$

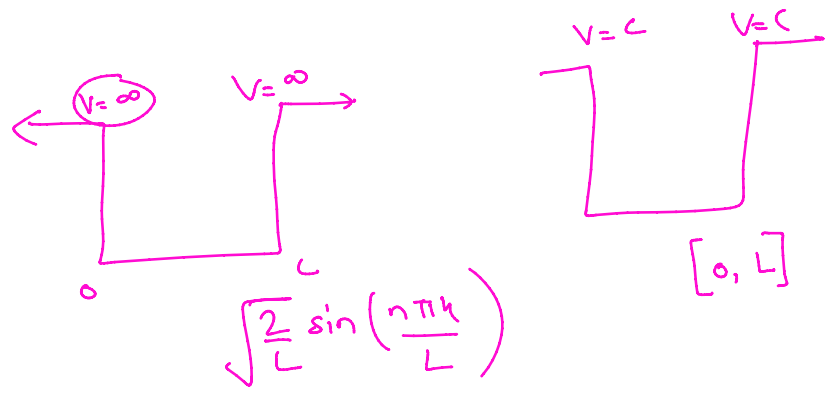
$$\Psi(x, t) = A e^{i(kx - \omega t)}$$

5. A particle moving in one dimension is in a stationary state whose wave function,

$$\Psi(x) = \begin{cases} 0, & x < -a \\ A(1 + \cos \frac{\pi x}{a}), & -a \leq x \leq a \\ 0, & x > a \end{cases}$$

where  $A$  and  $a$  are real constants.

- (a) Is this a physically acceptable wave function? Explain.  
 (b) Find the magnitude of  $A$  so that  $\psi(x)$  is normalized.  
 (c) Evaluate  $\Delta x$  and  $\Delta p$ . Verify that  $\Delta x \Delta p \geq \hbar/2$ . ( $\hbar/2$ )  
 (d) Find the classically allowed region.



So? (i)  $\varphi(-a^-) = 0$   $\varphi(-a^+) = 0$   
 $\varphi(a^-) = 0$   $\varphi(a^+) = 0$

(ii)  $\varphi'(a^-) = 0$   $\varphi'(a^+) = 0$   
 $\varphi'(-a^-) = 0$   $\varphi'(-a^+) = 0$

(iii) Normalizability

$$\int_{-\infty}^{\infty} \varphi^* \cdot \varphi \cdot dn = \int_{-a}^a |A|^2 \left(1 + \cos \frac{\pi n}{a}\right)^2 \cdot dn$$

$$= \int_{-a}^a |A|^2 \left[4 \cos^4 \frac{\pi n}{2a}\right] \cdot dn$$

$$= \frac{8a}{\pi} |A|^2 \int_{-\pi/2}^{\pi/2} \cos^4 t \cdot dt$$

$\frac{\pi n}{2a} = t$   $dn = \frac{2a \cdot dt}{\pi}$

$\downarrow$   
 $2 \times 3\pi/16$

$$= 3a |A|^2 \rightarrow \text{finite}$$

Hence  $\varphi(n, t)$  is normalizable

(b)  $3a |A|^2 = 1$

$$A = \frac{1}{\sqrt{3a}}$$

(c)  $\Delta n = \sqrt{\langle n^2 \rangle - \langle n \rangle^2}$

$$\langle n \rangle = \int_{-\infty}^{\infty} \varphi^*(n) \cdot n \cdot \varphi(n) \cdot dn = \int_{-a}^a \frac{4n}{3a} \cos^4 \left(\frac{\pi n}{2a}\right) \cdot dn = 0 \quad (\text{Why?})$$

$$\langle n^2 \rangle = \int_{-\infty}^{\infty} \varphi^*(n) \cdot n^2 \cdot \varphi(n) \cdot dn = \int_{-a}^a \frac{4n^2}{3a} \cos^4 \left(\frac{\pi n}{2a}\right) \cdot dn = \int_0^a \frac{8n^2}{3a} \cos^4 \left(\frac{\pi n}{2a}\right) \cdot dn$$

$$= \frac{8}{3a} \int_0^a n^2 \left[1 + \cos \frac{\pi n}{a}\right]^2 \cdot dn$$

$$= \frac{8}{3a} \int_0^a n^2 + 2n^2 \cos \frac{\pi n}{a} + n^2 \cos^2 \frac{\pi n}{a} \cdot dn$$

$$= \frac{8}{3a} \int_0^a \frac{3n^2}{2} + 2n^2 \cos \frac{\pi n}{a} + \frac{n^2}{2} \cos \frac{2\pi n}{a} \cdot dn$$

$$= \frac{8}{3a} \left[ \frac{a^3}{2} + \frac{2n^2 \sin \pi n / a}{\pi / a} \Big|_0^a - \frac{4}{(\pi/a)} \int_0^a n \sin \left( \frac{\pi n}{a} \right) + \frac{n^2 \sin \left( \frac{2\pi n}{a} \right)}{2(2\pi/a)} \right. \\ \left. - \frac{2}{2(2\pi/a)} \int_0^a n \sin \left( \frac{2\pi n}{a} \right) \right]$$

$$= \frac{8}{3a} \left[ \frac{a^3}{2} - \frac{4}{(\pi/a)} \left[ \frac{n \cos(\pi n/a)}{\pi/a} \Big|_0^a + \frac{\sin(\pi n/a)}{(\pi/a)^2} \Big|_0^a \right] \right.$$

$$\left. - \frac{1}{(2\pi/a)} \left[ \frac{n \cos(2\pi n/a)}{2\pi/a} \Big|_0^a + \frac{\sin(2\pi n/a)}{(2\pi/a)^2} \Big|_0^a \right] \right]$$

$$= \frac{8}{3a} \left[ \frac{a^3}{2} - \frac{4}{(\pi/a)} \left[ \frac{a^2}{\pi} \right] - \frac{1}{(2\pi/a)} \left[ -\frac{a^2}{2\pi} \right] \right]$$

$$= \frac{8}{3a} \left[ \frac{a^3}{2} - \frac{4a^3}{\pi^2} + \frac{a^3}{4\pi^2} \right] = \frac{8}{3a} \left[ \frac{a^3}{2} - \frac{15a^3}{4\pi^2} \right] = \frac{8a^2}{3} \left[ \frac{2\pi^2 - 15}{4\pi^2} \right]$$

$$\Delta n = \sqrt{\langle n^2 \rangle - \langle n \rangle^2} = \sqrt{\langle n^2 \rangle} = \sqrt{\frac{2a^2}{3} \left[ \frac{2\pi^2 - 15}{\pi^2} \right]}$$

$$\langle p \rangle = \int_{-a}^a \psi^*(n) \hat{p} \psi(n) \cdot dn = \int_{-a}^a A^* (1 + \cos \frac{\pi n}{a}) -i\hbar \frac{\partial}{\partial n} \left[ A (1 + \cos \frac{\pi n}{a}) \right] \cdot dn$$

$$= \int_{-a}^a \frac{i\hbar}{3a} \cdot \left( \frac{\pi}{a} \right) \left[ \sin \left( \frac{\pi n}{a} \right) + \frac{\sin(2\pi n/a)}{2} \right] \cdot dn$$

$$= \frac{i\hbar \pi}{3a^2} \cdot 0 = 0$$

$$\langle p^2 \rangle = \int_{-a}^a A^* (1 + \cos \frac{\pi n}{a}) \cdot -\hbar^2 \frac{\partial^2}{\partial n^2} A (1 + \cos \frac{\pi n}{a}) \cdot dn$$

$$\langle p^2 \rangle = \int_{-a}^a A^* (1 + \cos \frac{\pi x}{a}) \cdot -\hbar^2 \frac{\partial^2}{\partial x^2} (1 + \cos \frac{\pi x}{a}) dx$$

$$= + \frac{\hbar^2}{3a} \left( \frac{\pi}{a} \right)^2 \int_{-a}^a \left( \cos \frac{\pi x}{a} + \cos^2 \frac{\pi x}{a} \right) dx$$

$\underbrace{\cos^2 \frac{\pi x}{a}}_{(1 + \cos 2\pi x/a)/2}$

$$= \frac{\hbar^2}{3a} \left( \frac{\pi}{a} \right)^2 \left[ 0 + \frac{1}{2} \times 2a + 0 \right] = \frac{\hbar^2 \pi^2}{3a^2}$$

$$\Delta p = \sqrt{\langle p^2 \rangle - \langle p \rangle^2} = \sqrt{\langle p^2 \rangle}$$

(I was proud to be a TA until I solved this question)

$$\Delta x \cdot \Delta p = \sqrt{\frac{2a^2}{3} \left[ \frac{2\pi^2 - 15}{\pi^2} \right]} \cdot \frac{\hbar^2 \pi^2}{3a^2} = \sqrt{\frac{2\hbar^2}{9} [2\pi^2 - 15]} = \hbar \sqrt{\frac{2\pi^2 - 15}{9}} \approx \hbar \sqrt{\frac{5}{9}} > \frac{\hbar}{2}$$

(d) Classically allowed region  $\rightarrow x \in (-a, a)$   
 $\hookrightarrow \psi(x)$  is non zero

$\{V < E\}$   $\{V > E\}$   
Classically Allowed energy

$\psi(x,t) = \psi(x) \cdot T(t)$   $\rightarrow$  TISE  $\left( -\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x,t)}{\partial x^2} + V(x) \cdot \psi(x,t) = i\hbar \frac{\partial \psi(x,t)}{\partial t} \right)$

$|\psi(x,t)|^2 = \psi(x,t) \cdot \psi^*(x,t)$   
 $= \left( \psi(x) \cdot e^{-iEt/\hbar} \right) \cdot \left( \psi^*(x) \cdot e^{iEt/\hbar} \right)$   
 $= |\psi(x)|^2 \rightarrow$  Stationary state

$-\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x)}{\partial x^2} + V(x) \psi(x) = E \psi(x)$   
 $i\hbar \frac{\partial T(t)}{\partial t} = ET(t)$   
 $T(t) = e^{-iEt/\hbar}$

$\rightarrow$  All separable solutions correspond to stationary states.

$$dP = \psi(x,t) \cdot \psi^*(x,t) dx = \psi(x) \cdot \psi^*(x) (e^{-i\omega t} \cdot e^{i\omega t}) \cdot dx = |\psi(x)|^2 \cdot dx$$

$\downarrow$   
 Probability independent of time  
(Stationary)

Q. Check if given wavefunctions are stationary wavefunctions

(i)  $\psi_1(x,t) = C_1 e^{-i\omega_0 t} + 5C_3 e^{-i\omega_0 t} = (C_1 + 5C_3) e^{-i\omega_0 t}$   
 $\underbrace{C_1 e^{-i\omega_0 t}}_{\hbar\omega_0}$

(ii)  $\psi_2(x,t) = C_2 e^{-i(3\omega_0)t} + C_4 e^{-i(5\omega_0)t}$   
 $(3\hbar\omega_0 \text{ or } 5\hbar\omega_0) \rightarrow$  wavefunction collapse

$\rightarrow$  not stationary states

$i\hbar \frac{\partial}{\partial t}$  (Sturm or Sturm)  $\rightarrow$  Wavefunction  $\rightarrow$  stationary states are definite energy states

6. \* Consider the 1-dimensional wave function of a particle of mass  $m$ , given by

$\psi(x) = A \left(\frac{x}{x_0}\right)^n e^{-x/x_0} \rightarrow$  Hydrogen like atom

where,  $A$ ,  $n$  and  $x_0$  are real constants.

(a) Find the potential  $V(x)$  for which  $\psi(x)$  is a stationary state (It is known that  $V(x) \rightarrow 0$  as  $x \rightarrow \infty$ ).

(b) What is the energy of the particle in the state  $\psi(x)$ ?

Sol<sup>n</sup> 
$$-\frac{\hbar^2}{2m} \frac{d^2\psi(x)}{dx^2} + V(x)\psi(x) = E\psi(x)$$

$\downarrow$   
must be an energy eigenfunction

$$-\frac{\hbar^2}{2m} A \left[ \frac{n \cdot n^{n-1}}{x_0^n} e^{-x/x_0} + \frac{x^n}{x_0^n} e^{-x/x_0} \cdot \left(-\frac{1}{x_0}\right) \right] + V(x) \cdot A \left(\frac{x}{x_0}\right)^n e^{-x/x_0} = E\psi(x)$$

$$-\frac{\hbar^2}{2m} A \left[ \frac{n(n-1)x^{n-2}}{x_0^n} e^{-x/x_0} - \frac{nx^{n-1}}{x_0^{n+1}} e^{-x/x_0} + \frac{nx^{n-1}}{x_0^n} e^{-x/x_0} \left(-\frac{1}{x_0}\right) - \frac{x^n}{x_0^{n+2}} e^{-x/x_0} \right] + V(x) \psi(x)$$

$$A \left(\frac{x}{x_0}\right)^n e^{-x/x_0} \left[ \frac{\hbar^2}{2m} \left[ \frac{n(n-1)}{x_0^2} - \frac{2n}{x_0} - \frac{1}{x_0^2} \right] + V(x) \right] = E\psi(x)$$

$$\Rightarrow V(x) = \frac{\hbar^2}{2m} \left[ \frac{n(n-1)}{x_0^2} - \frac{2n}{x_0} \right]$$

$V(x) \rightarrow 0$  as  $x \rightarrow \infty$   
 $C = 0$

(b) 
$$E = \frac{\hbar^2}{2m x_0^2}$$