

Tutorial 3 Solution

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1. * Consider a vector field $\vec{F}(r)$, where $r = \vec{r}$ and $\vec{F}(r)$ dies faster than $\frac{1}{r}$ as $r \rightarrow \infty$, show the following results

(a) Using Helmholtz theorem as discussed in Lecture 5, Show that $\vec{F}(r)$ may be written as

$$\vec{F}(r) = -\nabla \int_V \frac{\nabla' \cdot \vec{F}(r')}{4\pi |r-r'|} dr' + \nabla \times \int_V \frac{\nabla' \times \vec{F}(r')}{4\pi |r-r'|} dr'$$

(b) Derive the same expression for $\vec{F}(r)$ using

$$\vec{F}(r) = \int_V d\tau' \vec{F}(r') \delta^3(r-r')$$

boundary of the integral is to be understood at ∞ .

Hint: Use the following

- (i) $-4\pi \delta^3(r-r') = \nabla^2 \frac{1}{|r-r'|}$
- (ii) $\nabla \times \nabla \times = \nabla \nabla \cdot - \nabla^2$
- (iii) $\nabla \frac{1}{|r-r'|} = -\nabla' \frac{1}{|r-r'|}$
- (iv) $\nabla \times \frac{\vec{F}(r')}{|r-r'|} = -\vec{F}(r') \times \nabla \left(\frac{1}{|r-r'|} \right)$ and 7(b) from Problem Set 2.

Soln

$$(b) \vec{F}(r) = \int_V d\tau' \vec{F}(r') \left(-\frac{1}{4\pi} \right) \nabla^2 \frac{1}{|r-r'|}$$

Position of the point where field is required \leftarrow location of source

$$\vec{A}(r) = \frac{1}{4\pi} \int_V \frac{\vec{C}(r') \cdot d\tau'}{|r-r'|}$$

$$= \frac{-1}{4\pi} \nabla^2 \int_V d\tau' \frac{\vec{F}(r')}{|r-r'|}$$

$$= \frac{1}{4\pi} \left[\nabla \times \nabla \times \int_V d\tau' \frac{\vec{F}(r')}{|r-r'|} - \nabla \int_V d\tau' \frac{\nabla' \cdot \vec{F}(r')}{|r-r'|} \right]$$

$$= \frac{1}{4\pi} \left[-\nabla \times \int_V d\tau' \left(\nabla' \frac{1}{|r-r'|} \right) \times \vec{F}(r') \right] - \nabla \int_V d\tau' \frac{\vec{F}(r') \cdot \nabla \left(\frac{1}{|r-r'|} \right)}{|r-r'|}$$

$$= \frac{1}{4\pi} \left[-\nabla \times \int_V d\tau' \left[\vec{F}(r') \times \nabla' \left(\frac{1}{|r-r'|} \right) \right] + \nabla \int_V \left[\vec{F}(r') \cdot \nabla' \left(\frac{1}{|r-r'|} \right) \right] \right]$$

$$= \frac{1}{4\pi} \left[-\nabla \times \int_V d\tau' \left[\vec{F}(r') \times \left(\frac{\vec{r}-\vec{r}'}{|\vec{r}-\vec{r}'|^3} \right) \right] + \nabla \int_V \left[\vec{F}(r') \cdot \left(\frac{\vec{r}-\vec{r}'}{|\vec{r}-\vec{r}'|^3} \right) \right] \right]$$

$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

$$\nabla \times \vec{E} = 0$$

$$\vec{E}(\vec{r}) = E_x(\vec{r})\hat{x} + E_y(\vec{r})\hat{y} + E_z(\vec{r})\hat{z}$$

$$\nabla \cdot \vec{E}(\vec{r}) = \rho(\vec{r})$$

$$\nabla \times \vec{E}(\vec{r}) = \vec{C}(\vec{r})$$

$\rho(\vec{r})$ & $\vec{C}(\vec{r})$ should fall off at

$$\vec{E}(\vec{r}) = -\nabla V(\vec{r}) + \nabla \times \vec{A}(\vec{r})$$

Scalar & rate faster than $\frac{1}{r^2}$

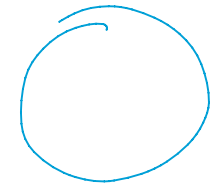
$$V(\vec{r}) = \frac{1}{4\pi} \int_V \frac{\rho(\vec{r}') \cdot d\tau'}{|\vec{r}-\vec{r}'|}$$

Position of the point where field is required \leftarrow location of source

$$\vec{A}(\vec{r}) = \frac{1}{4\pi} \int_V \frac{\vec{C}(\vec{r}') \cdot d\tau'}{|\vec{r}-\vec{r}'|}$$

$$= \frac{1}{4\pi} \left[-\nabla \times \int \frac{d\vec{c}}{|\vec{r}-\vec{r}'|} + \nabla \left(\frac{\nabla \cdot \vec{F}(\vec{r}')}{|\vec{r}-\vec{r}'|} \right) \right]$$

$$= \frac{1}{4\pi} \left[-\nabla \times \int \frac{\vec{F}(\vec{r}') \times d\vec{S}'}{|\vec{r}-\vec{r}'|} + \nabla \left(\int \frac{\nabla \cdot \vec{F}(\vec{r}')}{|\vec{r}-\vec{r}'|} d\tau' \right) \right]$$



$\nabla f =$ $R = \infty$

$$= \left[\frac{1}{|\vec{r}-\vec{r}'|} \right] + \frac{1}{4\pi} \left[\int \frac{\vec{F}(\vec{r}') \cdot d\vec{S}'}{|\vec{r}-\vec{r}'|^2} - \int \frac{\nabla \cdot \vec{F}(\vec{r}')}{|\vec{r}-\vec{r}'|} d\tau' \right]$$

$$\nabla \int \frac{\vec{F}(\vec{r}') \cdot d\vec{S}'}{|\vec{r}-\vec{r}'|}$$

$\vec{r} = \vec{r} - \vec{r}'$

$\frac{1}{|\vec{r}-\vec{r}'|^2} \hat{r}$

$\int \frac{\vec{F}(\vec{r}') \cdot \hat{r}}{r^2} d\vec{S}'$

$\int \frac{\vec{F}(\vec{r}') \cdot \hat{r}}{r^2} d\vec{S}' = \int \frac{\vec{F}(\vec{r}') \cdot \hat{r}}{r^2} r^2 \sin\theta d\theta d\phi$

Soln

$$\lim_{r \rightarrow \infty} \nabla \left[\int \frac{\vec{F}(\vec{r}') \cdot d\vec{S}'}{|\vec{r}-\vec{r}'|} \right]$$

$$\int \frac{\vec{F}(\vec{r}') \cdot \hat{r}}{r^2} r^2 \sin\theta d\theta d\phi$$

$$\nabla \int \frac{\vec{F}(\vec{r}') \cdot \hat{r}}{r^2} r^2 \sin\theta d\theta d\phi$$

$$\frac{\vec{F}(\vec{r}') \cdot \hat{r}}{|\vec{r}-\vec{r}'|^2}$$

4. Evaluate the following integral

$$\int_V \mathbf{r} \cdot (\mathbf{d} - \mathbf{r}) \delta^3(\mathbf{e} - \mathbf{r}) d\tau$$

where $\mathbf{d} = (5, 5, 5)$, $\mathbf{e} = (15, 19, 17)$, and V is a sphere of radius 7 centered at $(10, 15, 19)$.

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* After an extremely precise measurement, it was revealed that the actual force between two point charges is given by -

$$\mathbf{F} = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r^2} \left(1 + \frac{r}{\lambda} \right) e^{-r/\lambda} \hat{r}$$

Where λ is a constant with dimensions of length, and it is a huge number which is why the correction is tiny and difficult to notice. Does this electric field results from a scalar potential? Justify. And if yes, find the potential due to a point charge q placed at the origin using infinity as your reference.

Solⁿ $\nabla \times \underbrace{\left(\frac{q_2}{4\pi\epsilon_0 r^2} \left(1 + \frac{r}{a}\right) e^{-r/a} \cdot \hat{r} \right)}_{\text{Field in radial}} = 0$

$\text{div}^n \rightarrow$ generated from a scalar pot.

$\nabla \phi = \left(\right)$

$\frac{\partial \phi}{\partial r} = \frac{q_2}{4\pi\epsilon_0 r^2} \left(1 + \frac{r}{a}\right) e^{-r/a}$ $\frac{\partial \phi}{\partial \theta} = 0$ $\frac{\partial \phi}{\partial \phi} = 0$

② $\delta [g(u)] = \sum_m \frac{1}{|g'(u_m)|} \delta(u - u_m)$

$\int_{-\infty}^{\infty} f(u) \cdot \delta [g(u)] = \sum_m \int_{u_m - \epsilon_0}^{u_m + \epsilon_0} f(u) \cdot \delta [g(u)]$

$g(u_m) = 0 \quad \forall$
 u_1, u_2, \dots, u_m
 $g'(u_i) \neq 0 \quad \forall i \in \{1, 2, \dots, m\}$

$= \sum_m \int_{u_m - \epsilon_0}^{u_m + \epsilon_0} f(u) \cdot \delta \left[g(u_m) + g'(u_m) \cdot (u - u_m) + \frac{g''(u_m)}{2} (u - u_m)^2 + \dots \right]$

$\delta(au) = \frac{1}{|a|} \delta(u)$

$= \sum_m \int_{u_m - \epsilon_0}^{u_m + \epsilon_0} f(u) \cdot \delta [g'(u_m) (u - u_m)]$

$= \sum_m \int_{u_m - \epsilon_0}^{u_m + \epsilon_0} f(u) \cdot \frac{\delta(u - u_m)}{|g'(u_m)|}$

$= \frac{1}{|g'(u_m)|} \sum_m \int_{u_m - \epsilon_0}^{u_m + \epsilon_0} f(u) \cdot \delta(u - u_m)$

$\int_{-\infty}^{\infty} f(u) \cdot \delta g(u)$

$= \sum_m \left[\frac{1}{|g'(u_m)|} \int_{u_m - \epsilon_0}^{u_m + \epsilon_0} f(u) \cdot \delta(u - u_m) \right]$

$= \int_{-\infty}^{\infty} f(u) \left[\sum_m \frac{\delta(u - u_m)}{|g'(u_m)|} \right]$

$$\delta[g(n)] = \sum_m \frac{\delta(n-n_m)}{|g'(n_m)|}$$

